

Self-consistent field approximation

partition function from Gaussian transformation

in Ginzburg-Landau formulation (close to T_c) $\bar{\phi} = Jzm$

$$Z = \int \mathcal{D}m(\vec{r}) e^{-\beta F(m; T, H)} \quad \text{functional integral}$$

$$F(m; T, H) = F_0(H, T) + \int d^d r \left\{ \frac{A}{2} m(\vec{r})^2 + \frac{B}{4} m(\vec{r})^4 - H(\vec{r}) m(\vec{r}) + \frac{\kappa}{2} [\vec{\nabla} m(\vec{r})]^2 \right\}$$

expansion around $\langle m(\vec{r}) \rangle = 0$ $T > T_c$ $F_L = 0$

$$F = F_0 + F_L + \frac{1}{2} \int d^d r d^d r' m(\vec{r}) \left\langle \frac{\delta^2 F}{\delta m(\vec{r}) \delta m(\vec{r}')} \right\rangle m(\vec{r}')$$

Gaussian approximation

$$\left\langle \frac{\delta^2 F}{\delta m(\vec{r}) \delta m(\vec{r}')} \right\rangle = \left\{ A + 3B \langle m(\vec{r})^2 \rangle - \kappa \vec{\nabla}^2 \right\} \delta(\vec{r} - \vec{r}')$$

Self-consistent field approximation

$$\left\langle \frac{\delta^2 F}{\delta m(\vec{r}) \delta m(\vec{r}')} \right\rangle = \left\{ A + 3B \langle m(\vec{r})^2 \rangle - \kappa \vec{\nabla}^2 \right\} \delta(\vec{r} - \vec{r}')$$

Fourier transformation $m(\vec{r}) = \frac{1}{\sqrt{L^d}} \sum_{\vec{q}} m_{\vec{q}} e^{i\vec{q} \cdot \vec{r}}$ $m_{\vec{q}}^* = m_{-\vec{q}}$
 $m_{\pm \vec{q}} = m'_{\vec{q}} \pm i m''_{\vec{q}}$

$$F = F_0 + \frac{1}{2} \sum_{\vec{q}} \{ A + 3B \langle m^2 \rangle + \kappa q^2 \} m_{\vec{q}} m_{-\vec{q}} = F_0 + \sum_{\vec{q}} G^{-1}(\vec{q}) m_{\vec{q}} m_{-\vec{q}}$$

$$Z' = Z_0 \prod_{\vec{q}} \int dm_{\vec{q}} dm_{-\vec{q}} \exp \{ -\beta G^{-1}(\vec{q}) m_{\vec{q}} m_{-\vec{q}} \}$$

$$= Z_0 \prod_{\vec{q}} \int dm'_{\vec{q}} dm''_{\vec{q}} \exp \left\{ -\beta G^{-1}(\vec{q}) (m'^2_{\vec{q}} + m''^2_{-\vec{q}}) \right\}$$

Self-consistent field approximation

$$F = F_0 + \frac{1}{2} \sum_{\vec{q}} \{A + 3B\langle m^2 \rangle + \kappa q^2\} m_{\vec{q}} m_{-\vec{q}} = F_0 + \sum_{\vec{q}} G^{-1}(\vec{q}) m_{\vec{q}} m_{-\vec{q}}$$

$$Z' = Z_0 \prod_{\vec{q}} \int dm_{\vec{q}} dm_{-\vec{q}} \exp \{-\beta G^{-1}(\vec{q}) m_{\vec{q}} m_{-\vec{q}}\}$$

self-consistent field

$$\langle m^2 \rangle = \frac{1}{L^d} \int d^d r \langle m(\vec{r})^2 \rangle = \frac{1}{L^d} \sum_{\vec{q}} \langle m_{\vec{q}} m_{-\vec{q}} \rangle$$

$$= \frac{Z_0}{Z' L^d} \sum_{\vec{q}} \prod_{\vec{q}'} \int dm_{\vec{q}'} dm_{-\vec{q}'} m_{\vec{q}'} m_{-\vec{q}'} \exp \{-\beta G^{-1}(\vec{q}') m_{\vec{q}'} m_{-\vec{q}'}\}$$



$$\langle m^2 \rangle = \frac{k_B T}{L^d} \sum_{\vec{q}} G(\vec{q}) = \frac{1}{L^d} \sum_{\vec{q}} \frac{k_B T}{A + 3B\langle m^2 \rangle + \kappa q^2}$$

Self-consistent field approximation - susceptibility

$$\langle m^2 \rangle = \frac{k_B T}{L^d} \sum_{\vec{q}} G(\vec{q}) = \frac{1}{L^d} \sum_{\vec{q}} \frac{k_B T}{A + 3B \langle m^2 \rangle + \kappa q^2}$$

susceptibility

$$\begin{aligned} \chi(T) &= \beta \frac{1}{L^d} \int d^d r \, d^d r' \{ \langle m(\vec{r}) m(\vec{r}') \rangle - \langle m(\vec{r}) \rangle \langle m(\vec{r}') \rangle \} \\ &= \beta \langle m_{\vec{q}=0}^2 \rangle = G(\vec{q}=0) = \frac{1}{A + 3B \langle m^2 \rangle} \quad T > T_c \end{aligned}$$



$$\chi^{-1} = [A + 3B \langle m^2 \rangle] = A + \frac{3B k_B T}{L^d} \sum_{\vec{q}} \frac{1}{\chi^{-1} + \kappa q^2}$$

Self-consistent field approximation - instability

$$\chi^{-1} = [A + 3B\langle m^2 \rangle] = A + \frac{3Bk_B T}{L^d} \sum_{\vec{q}} \frac{1}{\chi^{-1} + \kappa q^2}$$

instability at $T \rightarrow T_c^*$ $\chi \rightarrow \infty \rightarrow \chi^{-1} \rightarrow 0$

$$A = -\frac{3Bk_B T}{L^d} \sum_{\vec{q}} \frac{1}{\kappa q^2} \rightarrow \frac{1}{L^d} \sum_{\vec{q}} \rightarrow \int \frac{d^d q}{(2\pi)^d} \rightarrow \frac{C_d}{(2\pi)^d} \int_0^\Lambda dq \frac{q^{d-1}}{\kappa q^2}$$

cutoff: $\Lambda \sim a^{-1}$

$$\begin{aligned} A_c &= Jz a^{-d} \left(\frac{T_c^*}{T_c} - 1 \right) \\ &= -\frac{Jz a^{-d}}{s^2} \frac{C_d k_B T_c^*}{(2\pi)^d} \int_0^\Lambda dq \frac{q^{d-1}}{\kappa q^2} \end{aligned}$$

$$T_c^* = \frac{T_c}{1 + \frac{C_d z}{(2\pi)^d} \frac{(\Lambda a)^{d-2}}{d-2}}$$

Self-consistent field approximation - critical exponents

$$\chi^{-1} = A + \frac{3Bk_BTC_d}{(2\pi)^d} \int_0^\Lambda dq \frac{q^{d-1}}{\chi^{-1} + \kappa q^2} \quad \longleftrightarrow \quad 0 = A_c + \frac{3Bk_BT_c^*C_d}{(2\pi)^d} \int_0^\Lambda dq \frac{q^{d-1}}{\kappa q^2}$$

temperature dependence of $\chi(T)$ for $T \rightarrow T_{c+}^*$

$$\chi^{-1} = (A - A_c) + \frac{3BC_d}{(2\pi)^d} \int_0^\Lambda dq q^{d-1} \left[\frac{k_BT}{\chi^{-1} + \kappa q^2} - \frac{k_BT_c^*}{\kappa q^2} \right]$$

$$\approx (A - A_c) - \frac{3BC_d k_B T_c^*}{(2\pi)^d \kappa} \int_0^\Lambda dq \frac{q^{d-3}}{1 + \chi \kappa q^2} \quad \leftarrow x = q\xi$$

$$= (A - A_c) - \frac{3BC_d k_B T_c^*}{(2\pi)^d \kappa} \underbrace{\{\kappa \chi\}^{(2-d)/2} \int_0^{\Lambda \sqrt{\kappa \chi}} dx \frac{x^{d-3}}{1+x^2}}_{\text{cutoff necessary for } d \geq d_c = 4}$$

correlation length $\xi^2 = \kappa \chi$

cutoff necessary for
 $d \geq d_c = 4$

Self-consistent field approximation - critical exponents

$$d \geq d_c = 4$$

$$\begin{aligned}\chi^{-1} &\approx (A - A_c) - \frac{3BC_d k_B T_c^*}{(2\pi)^d \kappa} \{\kappa\chi\}^{(2-d)/2} \frac{\{\Lambda(\kappa\chi)^{1/2}\}^{d-4}}{d-4} \\ &= \frac{k_B}{a^d s^2} \underbrace{(T - T_c^*)}_{\propto \tau} - \frac{C_d z^2}{2(2\pi)^d} \frac{T_c^*}{T_c} \frac{(\Lambda a)^{d-4}}{d-4} \chi^{-1}\end{aligned}$$

$$\Rightarrow \chi(T) = \frac{a^d s^2}{k_B (T - T_c^*)} \left\{ 1 + \frac{C_d z^2}{2(2\pi)^d} \frac{T_c^*}{T_c} \frac{(\Lambda a)^{d-4}}{d-4} \right\}^{-1} \propto |\tau|^{-1}$$

$$\begin{aligned}\chi &\propto |\tau|^{-\gamma} \\ \xi &\propto |\tau|^{-\nu}\end{aligned} \Rightarrow \left\{ \begin{array}{l} \xi^2 = \kappa\chi \\ \text{Fisher scaling} \\ \gamma = (2 - \eta)\nu \end{array} \right\} \Rightarrow$$

$$\begin{aligned}\gamma &= 1 && \text{mean field exponents} \\ \nu &= 1/2 && \text{but} \\ \eta &= 0 && T_c^* < T_{\text{cmf}}\end{aligned}$$

Self-consistent field approximation - critical exponents

$$d < d_c = 4$$

$$\frac{d-2}{2} < 1$$

$$\chi^{-1} = \underbrace{(A - A_c)}_{\propto \tau} - \frac{3BC_d k_B T_c^*}{(2\pi)^d \kappa^{d/2}} K_d \chi^{(2-d)/2}$$

regime 1: $\chi^{(2-d)/2}$ dominant T "close" to T_c^*

$$\chi \propto |\tau|^{-\gamma} \quad \gamma = \frac{2}{d-2}$$

regime 2: χ^{-1} dominant T "far" from T_c^*

$$\chi \propto |\tau|^{-\gamma} \quad \gamma = 1$$

mean field

